### Evolution of a Bubble in Extensional Flow

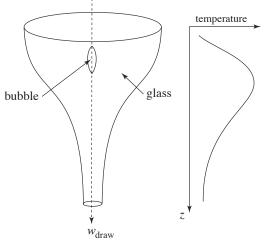
Academic Participants:

Elyse Fosse, Nicholas Gewecke, Alvaro Guevara, Saleem Ahmed, Kuan Xu, Burt Tilley, Linda Cummings, Don Schwendeman, Colin Please, Ellis Cumberbatch, Chris Breward, Giles Richardson

Industrial Presenter: John Abbott, Corning, Inc.

MPI 2008 at WPI

June 20, 2008



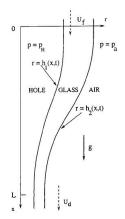
What shape is the final bubble?

- Fore/Aft symmetric?
- Same as uniform stretching?
- Longer/shorter than expected?
- What controls final bubble shape?
  - Temperature profile
  - Glass viscosity
  - Surface tension
  - Gas pressure
  - Initial bubble shape
- Are previous analyses of bubbles in shear flow relevant?

- Newtonian viscous flow
- Variable viscosity due to temperature (prescribed)
- Constant surface tension
- Axisymmetric flow
- Ideal gas with spatially uniform properties in the bubble

"The Mathematical Modelling of Capillary Drawing for Holey Fibre Manufacture," Fitt, et.al.

- Optical fibers that contain air holes, inner and outer radii of air hole is taken into consideration
- Gave three governing equations:
  - Conservation of mass
  - Normal stress condition on each surface
- Our model is obtained by making the inner radius much less than the outer radius



# Equations

### Glass Behavior

$$\frac{\partial}{\partial z} \left[ wR^2 \right] = 0$$

$$R^2 \operatorname{Re} \left[ w \frac{\partial w}{\partial z} \right] = \frac{\partial}{\partial z} \left[ 3R^2 \mu(T(z)) \frac{\partial w}{\partial z} + \sigma R \right]$$

### **Bubble Behavior**

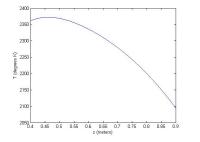
$$\frac{\partial h^2}{\partial t} + \frac{\partial}{\partial z} \left[ h^2 w \right] = \frac{1}{\mu(T(z))} \left[ h^2 P - \sigma h \right]$$
$$\frac{P}{R_g} \int_{z_{\min}}^{z_{\max}} \frac{\pi h^2(z,t)}{T(z)} dz = M.$$

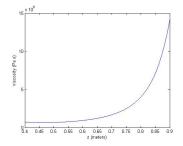
< < >> < <</>

Parameter		Range
Length		1 <i>m</i>
Glass Radius	R(z)	$10^{-1} - 10^{-4} m$
Bubble Radius	h(z,t)	10 <sup>-5</sup> – 10 <sup>-8</sup> m
Viscosity	$\mu(T(z))$	$10^5 - 10^6 \ kgm^{-1}s^{-1}$
Axial Velocity	W(Z)	$10^{-3} - 10^1 \ ms^{-1}$
Surface Tension	σ	.25 Nm <sup>-1</sup>

 $\langle \Box \rangle \langle \Box \rangle$ 

### Temperature and Viscosity profiles





### Question

At steady state, do the equations which govern the velocity *w* result in a velocity profile which fits the provided data?

We try to answer this question using numerical methods, by focusing on the first two equations.

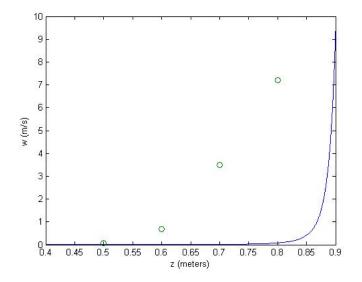
Assume Reynolds number,  $\sigma$  is small.

Equations  
$$\frac{\partial}{\partial z} \left[ wR^2 \right] = 0$$
$$\frac{\partial}{\partial z} \left[ 3R^2 \mu(T(z)) \frac{\partial w}{\partial z} \right] = 0$$

#### Method

- First equation  $\Rightarrow R^2 w = Q$ (constant)
- Set  $R^2 = Q/w \Rightarrow$  Second equation is DE in terms of w
- Integrate twice to get expression for w
- Apply BCs numerically and solve for *w* numerically

# **Basic approximation**



#### Neglect only inertia

Equations  

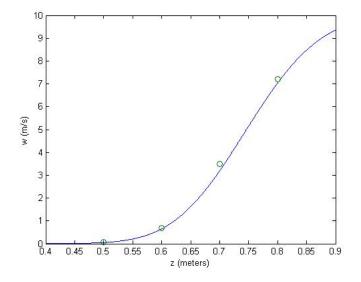
$$\frac{\partial}{\partial z} \left[ wR^2 \right] = 0$$

$$\frac{\partial}{\partial z} \left[ 3R^2 \mu(T(z)) \frac{\partial w}{\partial z} - \sigma R \right] = 0$$

#### Method

- Equation (1)  $\Rightarrow R^2 = Q$  (constant)
- Set  $R^2 = Q/w \Rightarrow$  Second equation is DE in terms of w
- Integrate once, then apply a shooting method on constant of integration (binary search and forward Euler) to find w

# Second approximation



Surface tension and inertia effects are included. Rewrite the second equation with  $R^2 = Q/w$ :

### Equations

$$\frac{\partial}{\partial z} \left[ w R^2 \right] = 0$$

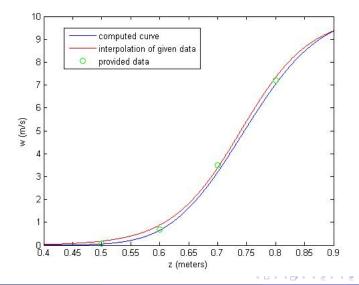
$$\frac{\partial}{\partial z} \left[ 3\frac{Q}{w} \mu(T(z)) \frac{\partial w}{\partial z} + \sigma \sqrt{\frac{Q}{w}} - Q \operatorname{Re} w \right] = 0$$

Follow same procedure as with Second approximation

Results are similar

# Comparison of Fits

A tanh fit to the data was used for other work on this problem



# **Bubble Dynamics**

Having found a good fit for w and R it is now possible to study the bubble behavior equations.

### **Governing Equations**

$$\frac{\partial h^2}{\partial t} + \frac{\partial}{\partial z} \left[ h^2 w \right] = \frac{1}{\mu(T(z))} \left[ h^2 P - \sigma h \right]$$
$$\int_{z_{\min}}^{z_{\max}} \rho_b(z, t) \pi h^2(z, t) dz = M, \qquad P(t) = \rho_b(z, t) R_g T(z)$$

- w(z) given from previous analysis
- initial conditions: h(z, 0) for  $z_{\min} \le z \le z_{\max}$  and P(0) given

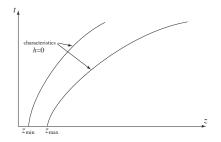
### Analysis

- Pinch-off behavior
- Numerical solution

Image: Image:

# Pinch-off Analysis

#### General solution behavior



- h(z, t) evolves on characteristics starting between  $z_{min}$  and  $z_{max}$
- If no pinch-off occurs, bubble stretches according to limiting characteristics
- Bubble may pinch off, so bubble may be shorter than expected

#### Assumptions

- P,  $\mu$  and  $\sigma$  taken to be constant
- Velocity profile taken to be  $w(z) = \gamma z$  for a fixed  $\gamma$ .

Equation to solve:

$$\frac{\partial h^2}{\partial t} + \gamma z \frac{\partial h^2}{\partial z} = \left(\frac{P}{\mu} - \gamma\right) h^2 - \frac{\sigma}{\mu} h$$

subject to h(z, 0) given. Solve using method of characteristics

## Solution

Along characteristics,

$$h(t) = \left[\frac{(\alpha h_0 - \beta) \exp\left(\frac{\alpha t}{2}\right) + \beta}{\alpha}\right],$$

while h > 0, otherwise h = 0, where

$$\alpha = \frac{P}{\mu} - \gamma, \qquad \beta = \frac{\sigma}{\mu}.$$

Result:

- For  $P > \gamma \mu$ , no pinch off occurs provided  $h_0 > \frac{\beta}{\alpha}$ .
- For  $P < \gamma \mu$ , the bubble pinches off at

$$\tilde{t} = \frac{2}{-\alpha} \ln \left( \frac{\beta + (-\alpha)h_0}{\beta} \right)$$

for any positive value of  $h_0$ .

## Numerical Solution of Bubble Dynamics

Characteristic form:

$$\frac{dh^2}{dt} = -w_z h^2 + \frac{1}{\mu(T(z))} \left[ Ph^2 - \sigma h \right], \quad \text{on } \frac{dz}{dt} = w(z)$$

Fixed mass constraint:

$$\frac{P}{R_g}\int_{z_{\min}}^{z_{\max}}\frac{\pi h^2(z,t)}{T(z)}\,dz=M$$

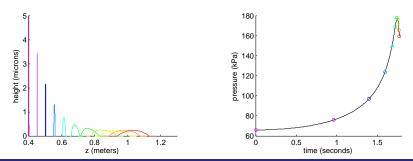
#### Basic numerical scheme (assuming w(z) is known):

- Pick mesh points  $\{z_j(0)\} \in [z_{\min}, z_{\max}]$ , specify  $h_j(0)$  (ICs).
- Advance  $z_j(t)$  and  $h_j(t)$  to  $t + \Delta t$  using the characteristic form with P(t) held fixed.
- Compute P(t + Δt) using fixed mass constraint.

## Results

Initial state:

- Ellipsoidal Bubble: 1mm length, 5µm radius (maximum)
- Bubble pressure balanced by surface tension

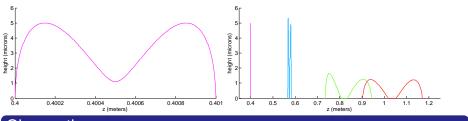


#### Observations

- Bubble radius shrinks due to stretching
- Smaller bubble radius requires higher pressure to balance surface tension
- Bubble shrinks more to raise pressure
- Pressure reduction occurs when stretching ceases and temperature continues to drop

Initial state:

- Dumb-bell shaped Bubble: 1mm length, 5μm radius (maximum), 1μm radius (minimum near center)
- Bubble pressure balanced by surface tension



### Observation

Bubbles may pinch off in the middle and break in two

Extensions to previous model:

- Gas may diffuse from bubble into surrounding glass (locally)
- Bubble mass no longer conserved
- Gas flux due to diffusion and bubble elongation

Model Equations:

$$\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial z} (w\rho) + \frac{\partial}{\partial r} (u\rho) = \frac{D}{r} \frac{\partial}{\partial r} (r \frac{\partial \rho}{\partial r})$$

where  $\rho(z, t)$  is density of gas in glass, h(z, t) is bubble radius, and  $u = \frac{h}{r} \frac{\partial h}{\partial z}$ .

Initial data for  $\rho$  specified, and the boundary condition  $\rho$  at r = h determined from pressure in bubble

Gas conservation of bubble ( $\rho_b$  = bubble density)

$$\int_{z_{\min}}^{z_{\max}} \rho_b \pi h^2 \, dz = M, \quad \frac{dM}{dt} = \int_{z_{\min}}^{z_{\max}} D \left. \frac{\partial \rho}{\partial r} \right|_{r=h} 2\pi h \, dz.$$

Pressure inside bubble related to density by gas law

$$P(t) = \rho_b(z, t) R_g T(z)$$

and  $\rho = \Lambda \rho_b$  on r = h,  $\Lambda$  is the solubility coefficient. This must be coupled to original three equations for *w*, *R*, *h*.

To be done!

## Summary & future directions

- Formulated a model for finite length axisymmetric bubble in glass fibre undergoing extensional flow (c.f. Fitt et al. "holey fibre" model)
- Four coupled equations for axial velocity w(z, t), fibre radius R(z, t), bubble radius h(z, t), and bubble pressure P(t). Furnace temperature T(z) (and hence viscosity) specified; mass of gas in bubble fixed.
- Model predicts that bubble cannot extend beyond limiting characteristics.
- Pinch-off model demonstrates that finite-time "pinch-off" of bubbles can occur, causing shorter bubbles
- Numerical code developed to solve governing equations for a fixed mass bubble
- Further numerical investigation of model is needed to investigate other possible phenomena under different drawing conditions and with different parameter values
- Formulated a model in which gas may leak from the bubble into surrounding glass – this model has yet to be studied